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# Classification of Lagrangian stars and their symplectic reductions 

S Janeczko $\dagger$<br>Institute of Mathematics, Warsaw University of Technology, Pl. Politechniki 1, 00-661, Warsaw, Poland

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#### Abstract

The Lagrangian star is a germ of the system ( $\left\{L_{1}, \ldots, L_{k}\right\}, \tilde{p}$ ) of Lagrangian submanifolds in the symplectic manifold $(M, \omega)$. We investigate the symplectic group action on Lagrangian stars and construct the basic invariants of such action. The Kashiwara signature for 3-Lagrangian linear stars is generalized to the nonlinear case and the generalized contact classes for Lagrangian stars are constructed. Finally, we obtain the generic classification of simple normal forms of reduced Lagrangian stars with respect to a hypersurface.


## 1. Introduction

Let $l_{1}, l_{2}, l_{3}$ be three Lagrangian subspaces in the symplectic vector space $(M, \omega)$. The natural invariant of the group of symplectic transformations of $M$, acting on the triplets of Lagrangian subspaces, is a signature (Maslov index [6]) of the Kashiwara quadratic form $Q\left(x_{1}, x_{2}, x_{3}\right)=\omega\left(x_{1}, x_{2}\right)+\omega\left(x_{2}, x_{3}\right)+\omega\left(x_{3}, x_{1}\right)$ defined on the direct sum $l_{1} \oplus l_{2} \oplus l_{3}$. In this paper we generalize this notion to the case of germs of triplets of Lagrangian submanifolds in a symplectic manifold. The problem considered is related to the classification of Lagrangian germs with respect to the subgroups of the group of symplectomorphisms. The natural subgroups are induced by $f$-liftable (cf [1]) vector fields $V$ on $M$ such that $\mathrm{d}(V\rfloor \omega)=0$ and $f$ is a smooth mapping between two manifolds, $f: N^{2 n} \rightarrow M^{2 n}$. Using the action of these groups one investigates the geometry of the maximal isotropic submanifolds in the degenerated symplectic structures (cf [3,7]) and show the direct way of generalizing the Lagrangian singularities. Using the symplectic invariants of contact (cf [4, 5]), in section 2 we find the algebraic invariants of the triplets of Lagrangian submanifolds containing two transversal submanifolds (basic Lagrangian star). We show that for the special class of tangential Lagrangian stars these invariants are determined by the equivalence class of right equivalence in the space of function-germs on $R^{n}$. Classification of reduced Lagrangian stars and basic Lagrangian stars, on a hypersurface $H$, under some genericity conditions is given in section 3. As an extension of this result the reduced local models, in the case of some non-transversal positions of Lagrangian stars with respect to $H$, are calculated.

## 2. Lagrangian stars

Let $(M, \omega)$ be a symplectic manifold. Let $\left\{L_{1}, \ldots, L_{k}\right\}$ be a system of Lagrangian submanifolds of ( $M, \omega$ ) intersecting at the common point $\tilde{p} \in L_{1} \cap \cdots \cap L_{k}$.
$\dagger$ E-mail address: janeczko@alpha.im.pw.edu.pl

Definition 2.1. The germ of Lagrangian submanifolds ( $\left\{L_{1}, \ldots, L_{k}\right\}, \tilde{p}$ ) is called a $k$-Lagrangian star at $\tilde{p}$. If $k=2$ and $L_{1}$ is transversal to $L_{2}$ then the 2-Lagrangian star ( $\left\{L_{1}, L_{2}\right\}, \tilde{p}$ ) is called the basic Lagrangian star. The 3-Lagrangian star we simply call the Lagrangian star.

Let $\left(\left\{L_{1}, \ldots, L_{k}\right\}, \tilde{p}\right)$ and $\left(\left\{L_{1}^{\prime}, \ldots, L_{k}^{\prime}\right\}, \tilde{p}\right)$ be two $k$-Lagrangian stars at $\tilde{p}$. Then we say that they are symplectically equivalent (or equivalent) if there is a germ of symplectomorphism $\Phi:((M, \omega), \tilde{p}) \rightarrow((M, \omega), \tilde{p})$ such that $\Phi\left(L_{j}\right)=L_{i_{j}}$ for some permutation $i_{j}$ of $\{1, \ldots, k\}$ and $\Phi(\tilde{p})=\tilde{p}$. The basic Lagrangian star forms a system of local symplectic coordinates of $(M, \omega)$. There are Darboux coordinates around $\tilde{p} \in M$ such that the basic Lagrangian $\operatorname{star}\left(\left\{L_{1}, L_{2}\right\}, \tilde{p}\right)$ is symplectically equivalent to the one defined by $L_{1}=\left\{(p, q) \in R^{2 n}, p=0\right\}$ and $L_{2}=\left\{(p, q) \in R^{2 n}, q=0\right\}$ with $(M, \omega) \cong\left(R^{2 n}, \sum_{i=1}^{n} \mathrm{~d} p_{i} \wedge \mathrm{~d} q_{i}\right)$.

To classify the Lagrangian stars we have to introduce the notion of contact equivalence and subsequently the symplectic contact equivalence. Let $X, L_{1}, L_{2}$ be equi-dimensional submanifolds of $M$ with $p \in X \cap L_{1} \cap L_{2}$. Then we say that $L_{1}$ and $L_{2}$ have the same contact with $X$ at $p$ if there is a germ of diffeomorphism $\phi:(M, p) \rightarrow(M, p)$ such that $\phi\left(L_{1}\right)=L_{2}$ and $\phi(X)=X$. Orbits of the group of these defined contact equivalences are called the contact classes. Using this definition, if $L_{1}, L_{2}$ have the same contact with $X$ at $p$ then the local rings $\bar{R}\left(X, L_{i}\right)=\mathcal{E}_{X} / \rho_{1}\left(X, L_{i}\right)$, where $\mathcal{E}_{X}$ denotes the local ring of smooth function-germs on $X$ at $p$ and $\rho_{1}\left(X, L_{i}\right)$ denotes the ideal of germs of functions on $M$ at $p$ which vanish to first order on $L_{i}$ restricted to $X$, are isomorphic. The corresponding isomorphism is induced by the pullback map $\phi^{\star}, \phi^{\star} f=f \circ \phi$, for $f \in \mathcal{E}_{X}$. The converse statement is true provided additionally $\operatorname{dim} \bar{R}\left(X, L_{i}\right)<\infty$.

The group of symplectomorphism-germs of $((M, \omega), \tilde{p})$ is a subgroup of the group of diffeomorphism-germs of $(M, \tilde{p})$ so that the contact data is a much more subtle invariant. If $X, L_{1}, L_{2}$ are Lagrangian submanifolds then the natural symplectic contact data is a pair

$$
\left(\bar{R}_{s}=\mathcal{E}_{X} / \rho_{2}\left(X, L_{i}\right), \sigma_{i}\right)
$$

where $\rho_{2}\left(X, L_{i}\right)$ denotes the ideal of germs of functions on $M$ at $\tilde{p}$ which vanish to second order on $L_{i}$ restricted to $X$, and the element $\sigma_{i} \in \bar{R}_{s}$ is naturally associated to $L_{i}$. In each case $\sigma$ is defined using a special cotangent bundle structure on a neighbourhood $\tilde{M}$ of $X$ in $M$, such that $\tilde{M}=T^{\star} X$ and $L_{i}=$ graph $\mathrm{d} \psi_{i}$ for some smooth functions $\psi_{i}$ on $X, \sigma_{i}$ is the image of $\psi_{i}$ in $\tilde{R}_{s}$. Obviously $\sigma$ is defined up to the choice of the special cotangent bundle structure $T^{\star} X$ on $M$. The special symplectic structure is a quadruple $(M, X, \pi, \theta)$, where $(M, X, \pi)$ is a differentiable fibre bundle, $\theta$ is a 1 -form on $M, \mathrm{~d} \theta=\omega$, such that there exists a diffeomorphism $\alpha: M \rightarrow T^{\star} X$ such that $\pi=\pi_{X} \circ \alpha, \theta=\alpha^{\star} \theta_{X}$. Let $L$ be a Lagrangian submanifold in $M$ and let $\theta_{1}$ and $\theta_{2}$ be 1-forms corresponding to two special symplectic structures on $(M, \omega)$ with the same base $X$. Then $\left.\theta_{1}\right|_{X}=\left.\theta_{2}\right|_{X}=0$ and near $X$ we have $\theta_{1}-\theta_{2}=\mathrm{d} H$, where $H$ is a function on $M$ which vanish to second order on $X$. The corresponding generating functions $\psi_{\theta_{1}}$ and $\psi_{\theta_{2}}$ of $L$ in both special symplectic structures are right equivalent with a diffeomorphism $g: X \rightarrow X$ defined by the formula

$$
g^{\star} \psi_{\theta_{2}}=\psi_{\theta_{1}}+\sum_{i j=1}^{n} h_{i j}\left(x, \mathrm{~d} \psi_{\theta_{1}}\right) \frac{\partial \psi_{\theta_{1}}}{\partial x_{i}} \frac{\partial \psi_{\theta_{1}}}{\partial x_{j}}
$$

where $H=\sum_{i j=1}^{n} h_{i j}(x, p) p_{i} p_{j}$. From this consideration we easily see the geometric sense of the local ring

$$
\bar{R}_{s}=\mathcal{E}_{X} /\left\langle\frac{\partial \psi}{\partial x_{1}}, \ldots, \frac{\partial \psi}{\partial x_{n}}\right\rangle^{2} .
$$

Now we assume that the Lagrangian star $S=\left(\left\{L_{1}, L_{2}, L_{3}\right\}, \tilde{p}\right)$ contains the basic Lagrangian star, say ( $\left.\left\{L_{1}, L_{2}\right\}, \tilde{p}\right)$. It is natural to define the pair of local rings $\boldsymbol{R}=\boldsymbol{R}_{1} \oplus \boldsymbol{R}_{2}$ associated to $S$, being a local invariant of the group of germs of symplectomorphisms acting on the space of Lagrangian stars (cf [4]). By considering germs of functions on $M$ near $\tilde{p}$ which vanish to second order on $L_{3}$ and taking the restrictions of these functions to $L_{1}$ (and respectively to $L_{2}$ ) we obtain an ideal $\Delta_{1}\left(L_{1}, L_{3}\right)$ (and respectively an ideal $\Delta_{2}\left(L_{2}, L_{3}\right)$ ).

Definition 2.2. By the basic invariant of the Lagrangian star $S$ we denote the pair of local rings

$$
\boldsymbol{R}=\boldsymbol{R}_{1} \oplus \boldsymbol{R}_{2}=\mathcal{E}_{L_{1}} / \Delta_{1}\left(L_{1}, L_{3}\right) \oplus \mathcal{E}_{L_{2}} / \Delta_{2}\left(L_{2}, L_{3}\right)
$$

where $\mathcal{E}_{L_{1}}$ (respectively $\mathcal{E}_{L_{2}}$ ) denotes the local ring of smooth function-germs on $L_{1}$ (respectively on $L_{2}$ ) near $\tilde{p}$. We call $S$ finite if $\operatorname{dim}_{R} \boldsymbol{R}<\infty$.

Now we have a natural realization of $\boldsymbol{R}_{1} \oplus \boldsymbol{R}_{2}$.
Proposition 2.1. For the considered Lagrangian star $S$

$$
\boldsymbol{R}_{i}=\mathcal{E}_{L_{i}} /\left\langle\frac{\partial \phi_{i}}{\partial x_{1}}, \ldots, \frac{\partial \phi_{i}}{\partial x_{n}}\right\rangle^{2}
$$

where $i=1,2$ and $\phi_{1}$ and $\phi_{2}$ are the function germs associated with the realizations of $L_{3}$ in two different cotangent bundle structures over $L_{1}$ and $L_{2}$. The generating functions $\phi_{1}$ and $\phi_{2}$ are defined up to an automorphism of $\boldsymbol{R}_{1}$ and $\boldsymbol{R}_{2}$ induced by the corresponding diffeomorphism-germs $L_{1} \rightarrow L_{1}$ and $L_{2} \rightarrow L_{2}$.

Proof. At first we recall some basic properties of Lagrangian submanifolds. If $X$ is a Lagrangian submanifold in $(M, \omega)$ then in some neighbourhood of $X$ the symplectic manifold $M$ is isomorphic to $T^{\star} X$. We say that $T^{\star} X$ is a special symplectic structure on $M$. Let $L$ be another Lagrangian submanifold in $(M, \omega)$, then around a point $p \in L \cap X$, the submanifold $L$ is generated by the generating function $F\left(p_{I}, q_{J}\right)$ (cf [1]), i.e. in local Darboux coordinates on $T^{\star} X, L$ is described by the equations

$$
\begin{equation*}
p_{J}=\frac{\partial F}{\partial q_{J}}\left(p_{I}, q_{J}\right) \quad q_{I}=-\frac{\partial F}{\partial p_{I}}\left(p_{I}, q_{J}\right) \tag{*}
\end{equation*}
$$

for some $J, I \subset\{1, \ldots, n\}, I \cap J=\emptyset, I \cup J=\{1, \ldots, n\}$. If the second equation of $(*)$ cannot be solved according to $p_{I}$ (around $\tilde{p}$ ) then obviously $L$ is vertical in directions $p_{I}$, so it cannot be generated by a function only on $q$. We see that $X$ is described by $\left\{p_{i}=0, i=1, \ldots, n\right\}$, so the ideal $\Delta(X, L)$ does not change if we perturb $L$ (make it transversal to the fibration $T^{\star} X \rightarrow X$ ) by adding the linear terms in $p$ to the second part of $(*)$ and making it solvable according to $p_{I}$. Thus we can represent the local ring $\boldsymbol{R}=\mathcal{E}_{X} / \Delta(X, L)$ by a generating function on $X$.

For the basic Lagrangian star $\left(\left\{L_{1}, L_{2}\right\}, \tilde{p}\right)$ we consider the special symplectic structures around $\tilde{p}, T^{\star} L_{1} \cong M$ and $T^{\star} L_{2} \cong M$. In both these structures the manifold $L_{3}$ can be defined by generating functions using the corresponding Liouville forms

$$
\left.\theta_{L_{i}}\right|_{L_{3}}=\mathrm{d} \tilde{\phi}_{i} \quad i=1,2
$$

We see that the ideals $\Delta_{i}\left(L_{i}, L_{3}\right)$ describe the order of contact of $L_{3}$ to $L_{i}$, (cf [4]) so by the small deformation of $L_{3}$ making it transversal to the fibrations $T^{\star} L_{i} \rightarrow L_{i}$ we get the generating functions $\phi_{i}$ of $L_{3}$ which may be defined on $L_{i}$ keeping $\Delta_{i}$ unchanged. These deformations may be achieved by changing the canonical 1-forms associated to the two cotangent bundle structures of $T^{\star} L_{i}$.

If the following two Lagrangian stars

$$
S=\left(\left\{L_{1}, L_{2}, L_{3}\right\}, \tilde{p}\right), S^{\prime}=\left(\left\{L_{1}^{\prime}, L_{2}^{\prime}, L_{3}^{\prime}\right\}, \tilde{p}\right)
$$

are symplectically equivalent then their corresponding basic invariants $\boldsymbol{R}_{1} \oplus \boldsymbol{R}_{2}$ and $\boldsymbol{R}_{1}^{\prime} \oplus \boldsymbol{R}_{2}^{\prime}$ are isomorphic. Now we would like to show that under certain conditions the converse is true.

Let the basic Lagrangian star of $S$ be in Darboux form, then $L_{3}$ is generated by a generating family $F_{q}\left(p_{I}, q\right)=S_{3}\left(p_{I}, q_{J}\right)+p_{I} q_{I}$ in $T^{\star} L_{1}$ and by $F_{p}\left(q_{J}, p\right)=S_{3}\left(p_{I}, q_{J}\right)-$ $p_{J} q_{J}$ in $T^{\star} L_{2}$, (which is the Legendre transform of $F_{q}$ ), for some $J, I \subset\{1, \ldots, n\}$, $I \cup J=\{1, \ldots, n\}$ and $I \cap J=\emptyset$. We choose $S_{3}$ such that

$$
\frac{\partial^{2} S_{3}}{\partial p_{I} \partial p_{I}}(0)=0
$$

This condition says that $L_{3}$ projects along $p$ with the kernel parametrized by $p_{I}$. In usual Lagrange equivalency preserving the fibration $(p, q) \rightarrow q$ we reduce $S_{3}$ to the form such that $S_{3}\left(p_{I}, q_{J}\right) \in \boldsymbol{m}_{I J}^{3}$. However, in this case we have to preserve the basic Lagrangian $\operatorname{star}\left(\left\{L_{1}, L_{2}\right\}, 0\right)$, where $L_{1}=\left\{(p, q) \in R^{2 n}, p=0\right\}$ and $L_{2}=\left\{(p, q) \in R^{2 n}, q=0\right\}$, so that the quadratic terms in some $q$-variables cannot be reduced. Thus we can write $S_{3}$ in the following final form

$$
S_{3}\left(p_{I}, q_{J}\right)=\tilde{S}\left(p_{I}, q_{J}\right)+Q\left(q_{J^{\prime}}\right)
$$

where $J^{\prime} \subset J, \tilde{S} \in \boldsymbol{m}_{I J}^{3}$ ( $\boldsymbol{m}_{I J}$ is the maximal ideal of smooth function-germs depending on $p_{I}, q_{J}$-variables, $\tilde{p}$ we assume to be 0 in these local coordinates), and $Q\left(q_{J^{\prime}}\right)$ is a nondegenerated quadratic form of $q_{J^{\prime}}$-variables, $\# J^{\prime}=l$. Now we can deduce the following result.

Proposition 2.2. Let $S$ and $S^{\prime}$ be two finite Lagrangian stars containing the basic Lagrangian star, then $S$ and $S^{\prime}$ are symplectically equivalent iff
(1) the quadratic forms $Q$ and $Q^{\prime}$ are equivalent, and
(2) the basic invariants $\boldsymbol{R}_{1} \oplus \boldsymbol{R}_{2}$ and $\boldsymbol{R}_{1}^{\prime} \oplus \boldsymbol{R}_{2}^{\prime}$ are isomorphic and the corresponding isomorphisms $\gamma_{1}$ and $\gamma_{2}$ send the images of $\phi_{1}$ and $\phi_{2}$ in $\boldsymbol{R}_{1}$ and $\boldsymbol{R}_{2}$ into the images of $\phi_{1}^{\prime}$ and $\phi_{2}^{\prime}$ in $\boldsymbol{R}_{1}^{\prime}$ and $\boldsymbol{R}_{2}^{\prime}$, respectively.

The basic invariant of the Lagrangian star $S$ is a $C^{\infty}$-invariant, i.e. an equivalence of Lagrangian stars is not necessarily symplectic. Now we see that the following data

$$
\left(\boldsymbol{R}, Q, \phi_{1}, \phi_{2}\right)
$$

form the complete symplectic invariant for Lagrangian stars under the symplectic group equivalence.

Remark 2.1. If $L_{3}$ is generated, in the basic Lagrangian star by $S_{3}\left(p_{I}, q_{J}\right), I \cup J=$ $\{1, \ldots, n\}, I \cap J=\emptyset$ and $\partial^{2} S_{3}(0) / \partial p_{I} \partial p_{I}=0$, then the class of $S$ is preserved if we apply the right equivalence group to $S_{3}$ preserving $\left\{p_{I}\right\}$ and $\left\{q_{J}\right\}$ spaces separately. There is a natural question, what does the relation between the local ring

$$
\mathcal{E}_{p_{I}, q_{J}} /\left(-\frac{\partial S_{3}}{\partial p_{I}}, \frac{\partial S_{3}}{\partial q_{J}}\right)^{2}
$$

and the basic symplectic invariant of the Lagrangian star $S$ look like?

Now we consider the special case. We assume $S$ contains the basic Lagrangian star. We call $S$ the tangential star if there are two Lagrangian germs in $S$ tangent at $\tilde{p}$. If $S$ is tangential then there exists local Darboux coordinates at $\tilde{p}$ in which $L_{1}=\left\{(p, q) \in R^{2 n}: p=0\right\}, L_{2}=\left\{(p, q) \in R^{2 n}: q=0\right\}$ and $L_{3}$ is generated by a generating function $q \rightarrow F(q), F^{\prime}(0)=0$ and $F^{\prime \prime}(0)=0$. The basic invariant for the tangential Lagrangian stars is reduced to the local ring

$$
\boldsymbol{R}=\mathcal{E}_{q} /\left\langle\frac{\partial F}{\partial q_{1}}, \ldots, \frac{\partial F}{\partial q_{n}}\right\rangle^{2}
$$

Symplectic equivalence of the tangential Lagrangian stars say $S$ and $S^{\prime}$, is equivalent to right equivalence of their generating functions $F$ and $F^{\prime}$. So the equivalence classes of contact are determined mainly by the $A_{k}, D_{k}$ and $E_{k}$ classification of singularities (cf [1]).

Remark 2.2. If $Q$ has a maximal rank then the main symplectic invariant of the Lagrangian star is a signature of $Q$. It is a signature (Maslov index) of the Kashiwara quadratic form (cf [6])

$$
\omega\left(x_{1}, x_{2}\right)+\omega\left(x_{2}, x_{3}\right)+\omega\left(x_{3}, x_{1}\right)
$$

defined on the tangent (at $\tilde{p}$ ) Lagrangian star $l_{1} \oplus l_{2} \oplus l_{3}$. We denote this signature by $\tau\left(l_{1}, l_{2}, l_{3}\right)$. This is a symplectic invariant for any Lagrangian star, not only if $l_{1} \cap l_{2}=l_{2} \cap l_{3}=l_{3} \cap l_{1}=\{0\}$. In general we can write

$$
\tau\left(l_{1}, l_{2}, l_{3}\right)=n+\operatorname{dim}\left(l_{1} \cap l_{2}\right)+\operatorname{dim}\left(l_{2} \cap l_{3}\right)+\operatorname{dim}\left(l_{3} \cap l_{1}\right)(\bmod 2)
$$

The basic symplectic invariant introduced here is a natural generalization of $\tau$ for the nonlinearizable Lagrangian stars. Generalization of this invariant for a Lagrangian star of four Lagrangian submanifolds goes through the composed 3-Lagrangian stars (cf [6]).

## 3. Reduction of Lagrangian stars

As far as the basic symplectic stars are all symplectically equivalent there is a natural question how they pass through the reduction on a hypersurface or a general co-isotropic submanifold? First, we consider the reduction of co-dimension 1, which is the very special reduction along the integral curves of the Hamiltonian system with the Hamiltonian function defining the hypersurface as its zero-level set.

Let $H$ be a hypersurface in $(M, \omega)$. We consider the basic stars $\left(\left\{L_{1}, L_{2}\right\}, \tilde{p}\right)$ transversal to $H$ at $\tilde{p}$. Let $\pi_{H}: H \rightarrow(\tilde{M}, \tilde{\omega})$ be the projection along bicharacteristics into the reduced symplectic manifold $\tilde{M}$ and $\tilde{\omega}$ is the corresponding reduced symplectic form $\pi_{H}^{\star} \tilde{\omega}=\left.\omega\right|_{H}$. We define the reduced star as follows

$$
\left(\left\{L_{1}^{r}, L_{2}^{r}\right\}, \tilde{p}^{r}\right)
$$

where $L_{i}^{r}=\pi_{H}\left(H \cap L_{i}\right)$ and $\tilde{p}^{r}=\pi_{H}(\tilde{p})$.
Now we pass to the classification of reduced basic stars according to the symplectomorphisms of $(M, \omega)$ and $(\tilde{M}, \tilde{\omega})$ preserving the projection $\pi_{H}$.

Proposition 3.1. Any simple, reduced basic Lagrangian star can be written in one form from the following normal forms

$$
A_{k}:\left(\left\{L_{1}^{r}, L_{2}^{r}\right\}, 0\right)
$$

where

$$
\begin{aligned}
L_{1}^{r} & =\left\{(q, p) \in \tilde{M}: p_{i}=0, i=1, \ldots, n-1\right\} \\
L_{2}^{r} & =\left\{(q, p) \in \tilde{M}: p_{1}=\partial S / \partial q_{1}\left(q_{1}, p_{2}, \ldots, p_{n-1}\right)\right. \\
& \left.q_{i}=-\partial S / \partial p_{i}\left(q_{1}, p_{2}, \ldots, p_{n-1}\right), i=2, \ldots, n-1\right\}
\end{aligned}
$$

and

$$
S\left(q_{1}, p_{2}, \ldots, p_{n-1}\right)= \pm q_{1}^{k+1}+q_{1}^{k-1} p_{k-1}+\cdots+q_{1}^{2} p_{2}
$$

for $2 \leqslant k \leqslant n-1$ and $A_{1}$ with $L_{2}^{r}=\left\{(q, p) \in \tilde{M}: q_{i}=0, i=1, \ldots, n-1\right\}$.
Proof. Now we have to classify the triplets $\left(\left\{L_{1}, L_{2}, H\right\}, \tilde{p}\right)$ in $\left(R^{2 n}, \omega\right)$, where $L_{1}, L_{2}$ and $H$ are mutually transversal at $\tilde{p} \in L_{1} \cap L_{2} \cap H$. We find Darboux coordinates $\left\{x_{1}, \ldots x_{n}, y_{1}, \ldots, y_{n}\right\}$ in which $L_{1}$ and $H$ may be written in the following normal form in $R^{2 n}=T^{\star} L_{1}$

$$
L_{1}=\left\{y_{1}=0, \ldots, y_{n}=0\right\} \quad H=\left\{x_{1}=0\right\}
$$

Then $L_{2}$ can be written by the generating function $y \rightarrow S(y)$, such that

$$
\mathrm{d}\left(\frac{\partial S}{\partial y_{1}}\right)(0) \neq 0
$$

because of transversality of $L_{1}$ to $L_{2}$ and $L_{2}$ to $H$. Now we need to use the symplectomorphisms of $T^{\star} L_{1}$ preserving $\left(L_{1} \cup H, 0\right)$ and reduce $L_{2}$ to its simple normal form. So we need the group $G_{L_{1} \cup H}$ of germs of symplectomorphisms which preserve the fibration $(x, y) \rightarrow y$ and the hypersurface $H=\left\{x_{1}=0\right\}$. Every element $\Phi$ of this group can be defined as a lifting of a diffeomorphism $\phi: R^{n} \ni y \rightarrow \phi(y) \in R^{n}$, which preserve the fibration over $\left(y_{2}, \ldots, y_{n}\right)$, i.e. $y \rightarrow \bar{y}=\left(y_{2}, \ldots, y_{n}\right)$ with adding the gradient of a function $f$ which depends on $\bar{y}$

$$
\Phi(x, y)=\left(\left(\phi^{\star}\right)^{-1}(y) x+\mathrm{d} f(\bar{y}), \phi(y)\right) .
$$

Using this group we can reduce the function $S$ to the form

$$
S(y)=y_{1} \tilde{S}(y)
$$

Then using the theorem on versal deformations (cf [8]) we reduce it successively to the form

$$
L_{2}: S(y)= \pm y_{1}^{k+1}+y_{1}^{k-1} y_{k}+\cdots+y_{1} y_{2} \quad \text { for } 1 \leqslant k \leqslant n
$$

At first we consider the case when $k=1$. In this case $S(y)= \pm y_{1}^{2}$ and the reduction equations $\left\{x_{1}=0, y_{1}=0\right\}$ give us the reduced Lagrangian germ in the form

$$
\pi_{\left\{x_{1}=0\right\}}\left(L_{2}\right)=\left\{(q, p) \in \tilde{M}: q_{i}=0, i=1, \ldots, n-1\right\}
$$

which corresponds to the case $A_{1}$ in the proposition.
Now we consider the case when $k \geqslant 2$. Taking the image $\pi_{\left\{x_{1}=0\right\}}\left(L_{2}\right)$ we obtain the following equations

$$
\begin{aligned}
& -x_{1}=\frac{\partial S}{\partial y_{1}}(y)= \pm(k+1) y_{1}^{k}+(k-1) y_{1}^{k-2} y_{k}+\cdots+2 y_{1} y_{3}+y_{2}=0 \\
& -x_{2}=\frac{\partial S}{\partial y_{2}}(y)=y_{1}, \ldots,-x_{k}=\frac{\partial S}{\partial y_{k}}(y)=y_{1}^{k-1} \\
& -x_{k+1}=\frac{\partial S}{\partial y_{k+1}}(y)=0, \ldots,-x_{n}=\frac{\partial S}{\partial y_{n}}(y)=0
\end{aligned}
$$

From the first equation we derive

$$
-y_{2}= \pm(k+1) y_{1}^{k}+(k-1) y_{1}^{k-2} y_{k}+\cdots+2 y_{1} y_{3}
$$

and renumerating the corresponding Darboux coordinates

$$
\begin{aligned}
& p_{1}=-y_{2}, p_{2}=y_{3}, \ldots, p_{n-1}=y_{n} \\
& q_{1}=-x_{2}, q_{2}=x_{3}, \ldots, q_{n-1}=x_{n}
\end{aligned}
$$

we rewrite the equations for $\pi_{\left\{x_{1}=0\right\}}\left(L_{2}\right)$ in the form

$$
\begin{aligned}
& p_{1}=\frac{\partial \bar{S}}{\partial q_{1}}= \pm(k+1) q_{1}^{k}+(k-1) q_{1}^{k-2} p_{k-1}+\cdots+2 q_{1} p_{2} \\
& q_{2}=-\frac{\partial \bar{S}}{\partial p_{2}}=-q_{1}^{2}, \ldots, q_{k-1}=-\frac{\partial \bar{S}}{\partial p_{k-1}}=-q_{1}^{k-1} \\
& q_{k}=-\frac{\partial \bar{S}}{\partial p_{k}}=0, \ldots, q_{n-1}=-\frac{\partial \bar{S}}{\partial p_{n-1}}=0
\end{aligned}
$$

with the generating function

$$
\bar{S}\left(q_{1}, p_{1}, \ldots, p_{n-1}\right)= \pm q_{1}^{k+1}+q_{1}^{k-1} p_{k-1}+\cdots+q_{1}^{2} p_{2}
$$

for the reduced Lagrangian germ $L_{2}^{r}$.
Remark 3.1. We see that the only stable case of the triplet $\left(\left\{L_{1}, L_{2}, H\right\}, \tilde{p}\right)$ is equivalent to the local model of type $A_{1}$ for the submanifold $L_{2}$ and that it corresponds to the basic reduced star which is the basic star in the reduced symplectic space.

Let $S=\left(\left\{L_{1}, L_{2}, L_{3}\right\}, \tilde{p}\right)$ be a Lagrangian star and let $H$ be a hypersurface-germ at $\tilde{p}$.
Proposition 3.2. We assume that the Lagrangian star $S$ contains the star, say ( $\left\{L_{1}, L_{2}\right\}, \tilde{p}$ ) which is of type $A_{1}$ (stable) with respect to ( $H, \tilde{p}$ ). Then in the transversal case i.e. $L_{3}$ is transversal to $L_{1}, L_{2}$, and $H$, the typical reduced stars $\left(\left\{L_{1}^{r}, L_{2}^{r}, L_{3}^{r}\right\}, \tilde{p}^{r}\right)$ are classified by the following normal forms: $\left(\left\{L_{1}^{r}, L_{2}^{r}\right\}, \tilde{p}^{r}\right)$ is a basic Lagrangian star in $\tilde{M} \equiv\left(R^{2(n-1)}, \omega=\sum_{i=1}^{n-1} \mathrm{~d} y_{i} \wedge \mathrm{~d} x_{i}\right)$ and $L_{3}^{r}$ is generated by the following Morse family

$$
F(\lambda, y)=\lambda^{k+1}+\sum_{i=1}^{k-1} \lambda^{k-i} y_{i}+\phi\left(y_{1}, \ldots, y_{k-1}\right) \pm y_{k}^{2} \pm \cdots \pm y_{n-1}^{2}
$$

where $\phi \in \boldsymbol{m}_{y_{1}, \ldots, y_{k-1}}^{2}$.
Proof. By proposition 3.1 we can reduce ( $\left\{L_{1}, L_{2}, H\right\}, \tilde{p}$ ) to the following normal form in $\tilde{p}$

$$
\begin{aligned}
& L_{1}: y_{1}=0, \ldots, y_{n}=0 \\
& L_{2}: x_{1}= \pm 2 y_{1}, x_{2}=0, \ldots, x_{n}=0 \\
& H: x_{1}=0
\end{aligned}
$$

By transversality assumptions $L_{3}$ can be generated by the generating function $y \rightarrow F(y)$ such that

$$
\mathrm{d}\left(\frac{\partial F}{\partial y_{1}}\right)(0) \neq 0
$$

By the reduction projection $\pi_{H}$ we get

$$
\begin{aligned}
& L_{1}^{r}: y_{2}=0, \ldots, y_{n}=0 \\
& L_{2}^{r}: x_{2}=0, \ldots, x_{n}=0 \\
& L_{3}^{r}: x_{2}=\frac{\partial F}{\partial y_{2}}\left(y_{1}, \bar{y}\right), \ldots, x_{n}=\frac{\partial F}{\partial y_{n}}\left(y_{1}, \bar{y}\right) \quad 0=\frac{\partial F}{\partial y_{1}}\left(y_{1}, \bar{y}\right)
\end{aligned}
$$

where $\bar{y}=\left(y_{2}, \ldots, y_{n}\right)$.
Any liftable (through $\pi_{H}$ ) equivalence of ( $\left\{L_{1}^{r}, L_{2}^{r}, L_{3}^{r}\right\}, 0$ ) is determined by an $\mathcal{R}$ equivalence of Morse families $F\left(y_{1}, \bar{y}\right)$, where the diffeomorphism of $\bar{y}$ is preserving zero. By reordering $\bar{y}$, treating $y_{1}$ as a Morse parameter, $\lambda$ and applying the group of equivalences we obtain the prenormal forms of proposition 3.2.

Now we consider the situation when the basic Lagrangian star ( $\left\{L_{1}, L_{2}\right\}, \tilde{p}$ ) is not transversal to the hypersurface $(H, \tilde{p})$. In this case at least one of the two Lagrangian germs $L_{1}, L_{2}$ have to be transversal to $H$. We assume that it is $L_{1}$. Then we have the following result.

Proposition 3.3. If the basic Lagrangian star $\left(\left\{L_{1}, L_{2}\right\}, \tilde{p}\right)$ is not transversal to ( $H, \tilde{p}$ ) then the generic reduced Lagrangian star $\left(\left\{L_{1}^{r}, L_{2}^{r}\right\}, \tilde{p}^{r}\right)$ can be written in one form from the following normal forms: $L_{1}^{r}: y_{1}=0, \ldots, y_{n-1}=0, L_{2}^{r}$ : is generated by the following Morse family

$$
S(\lambda, \bar{y})= \pm \lambda^{k}+\sum_{i=1}^{k-3} y_{i} \lambda^{k-i-1}+\left(g\left(y_{1}, \ldots, y_{k-3}\right) \pm \sum_{i=k-2}^{n} y_{i}^{2}\right) \lambda
$$

where $k=\operatorname{dim}_{R} \mathcal{E}_{\lambda} / \Delta(S(\lambda, 0))+1 \leqslant n+2, g \in \boldsymbol{m}_{y_{1}, \ldots, y_{k-3}}^{2}-\boldsymbol{m}_{y_{1}, \ldots, y_{k-3}}^{3}$ and $\Delta(S(\lambda, 0))$ is an ideal in $\mathcal{E}_{\lambda}$ generated by $\partial S / \partial \lambda(\lambda, 0)$.

Proof. If $L_{1}$ is transversal to $H$ then we obtain the Darboux coordinates such that

$$
L_{1}: y_{1}=0, \ldots, y_{n}=0 \quad H: x_{1}=0
$$

Because $L_{2}$ is transversal to $L_{1}$, then $L_{2}$ may be generated in the form

$$
x_{i}=-\frac{\partial S}{\partial y_{i}}(y) \quad i=1, \ldots, n
$$

where $\left.\mathrm{d}\left(\partial S / \partial y_{1}\right)\right|_{0}=0$. By the equivalence group of symplectomorphisms preserving ( $H \cup L_{1}, 0$ ) we can reduce $S$ to the form (cf [2])

$$
S(y)= \pm y_{1}^{k}+\sum_{i=2}^{k-2} y_{i} y_{1}^{k-i}+\left(g\left(y_{2}, \ldots, y_{k-2}\right) \pm \sum_{i=k-1}^{n} y_{i}^{2}\right) y_{1}
$$

where $g$ is a smooth function (functional invariant) and $g \in \boldsymbol{m}_{y_{2}, \ldots, y_{k-2}}^{2}$. By the reduction projection $\pi_{H}$ and reordering the variables ' $y$ ' we get the corresponding Morse family $S(\lambda, \bar{y}), \bar{y}=\left(y_{1}, \ldots, y_{n-1}\right)$, generating $L_{2}^{r}$.

We see that the reduction of the basic Lagrangian star, which is not transversal to $H$ is no more basic. Moreover it is not even smooth. The only simple model of the reduced Lagrangian star in the non-transversal case is the one generated by the following Morse family

$$
L_{2}^{r}: S\left(\lambda, y_{1}, y_{2}\right)=\lambda^{3} \pm y_{1}^{2} \lambda
$$

which is the singular Lagrangian set.

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